Improvement of modulation bandwidth in electroabsorption-modulated laser by utilizing the resonance property in bonding wire

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Abstract: We present and demonstrate a simple and cost-effective technique for improving the modulation bandwidth of electroabsorption-modulated laser (EML). This technique utilizes the RF resonance caused by the EML chip (i.e., junction capacitance) and bonding wire (i.e, wire inductance). We analyze the effects of the lengths of the bonding wires on the frequency responses of EML by using an equivalent circuit model. To verify this analysis, we package a lumped EML chip on the sub-mount and measure its frequency responses. The results show that, by using the proposed technique, we can increase the modulation bandwidth of EML from ~16 GHz to ~28 GHz.

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1. Introduction

Due to the ever-increasing data traffic in the fiber-optic network, there have been many efforts to develop high-speed light sources [1–6]. These light sources should also be compact,

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reliable, and cost-effective, and have superior performances such as the low chirp and low power consumption. It has been considered that the electro-absorption-modulated laser (EML) can satisfy all these requirements, thus is suitable for the use in the high-speed applications operating at 25 Gb/s or 40 Gb/s [3–6]. However, the modulation bandwidth of the EML is often limited by the RC time constant resulting from its junction capacitance. To overcome this problem, it is necessary to modify the structural parameters of the electro-absorption modulator (EAM). However, it is usually difficult to increase the modulation bandwidth without compromising other important parameters such as the extinction ratio, insertion loss, and required driving voltage.

There are several techniques commonly used to avoid these performance compromises [6–9]. For example, it has been demonstrated that the modulation bandwidth of the EML can be increased, without sacrificing other performances, by introducing the travelling-wave (TW) electrode to the EAM [6]. However, this technique requires complicated fabrication processes and stringent process tolerances. The modulation bandwidth of the EML can also be extended by utilizing the equalization technique based on the optical filter [7]. However, in spite of its excellent capability of enhancing the modulation bandwidth (e.g., from 10 to 45 GHz), this technique is problematic since it not only requires a large number of additional components and precise optical alignment, but also suffers from large excess loss.

Previously, it has also been reported that the modulation bandwidth of EAM can be improved by using the flip-chip bonding technique and controlling the inductance of the bonding lead between the EAM chip and load impedance [8]. However, when this flip-chip bonding technique applied to the EML, it can suffer from the difficulty of heat dissipation of distributed feedback laser diode (DFB LD) and the effect of RF parasitics of ceramic (or Si) sub-mount, and can require an additional alignment pattern for the chip (boned to the sub-mount in the form of p-side down). Another technique proposed to increase the modulation bandwidth of EML is using the RF resonance caused by placing a T-resonator (i.e., shunt open circuit stub) between the lumped EML chip and the load impedance on the sub-mount. Thus, the frequency response of the EML can be changed by the length of the stub [9]. This technique appears to be attractive since it can increase the modulation bandwidth of the EML without degrading other performances including the excess loss, and has a simple configuration. Previously, it has been shown that this technique can increase the modulation bandwidth of EML by only 2 GHz using a numerical simulation. No experimental data supporting this simulation result has yet been reported.

In this paper, we propose and demonstrate a simple and cost-effective technique for improving the modulation bandwidth of EML. This is based on the wire-bonding technique (widely used for most commercial transmitted optical subassembly (TOSA)) and utilizes the RF resonance generated by the junction capacitance of the EAM and the inductances of the bonding wires on its sub-mount. Compared to the flip-chip bonding technique, this technique offers sufficient heat-dissipation of the EML and has good immunity to the RF parasitics of sub-mount. Besides, there is no significant constraint to the geometry of the device and sub-mount (i.e., high degree of freedom) in doing the wire-bonding and no need to use expensive AuSn solder which is usually used for flip-chip bonding (i.e., cost-effective). The frequency response of the EML can be compensated simply by adjusting the lengths of bonding wires or by using the different types of bonding wire (i.e., simple configuration). To evaluate the effectiveness of the proposed technique, we package an EML chip on the sub-mount and investigate the effect of this inductance modification (caused by adjusting the lengths of the bonding wires between the RF feeding line and EML chip and also between the EML chip and load impedance) on its frequency responses.

The rest of this paper is organized as follows. The circuit model and simulation results are described in Section II; Section III presents the experimental results obtained by using the proposed technique. Finally, Section IV concludes the paper.

2. Circuit model and simulation results

Figure 1 shows the equivalent circuit diagram of EML. The CW light generated by the DFB LD (by injecting the source current I_s) is intensity-modulated by the temporal change of the optical absorption within the EAM (by modulating the source voltage V_s). From this circuit model, the relation between V_s and the voltage applied to the EAM (V_m) can be expressed as

$$V_{m} = \frac{Z_{w2} + Z_{L}}{R_{s} + Z_{w1} + Z_{w2} + Z_{L} + (R_{s} + Z_{w1})(Z_{2} + Z_{wL})/Z_{m}}V_{s}$$
(1)

where

$$Z_{m} = \frac{1}{\frac{1}{R_{m} + 1/\left(j\omega C_{m} + \frac{1}{R_{pho}} + \frac{1}{R_{j}}\right)} + j\omega C_{p}}$$
(2)

In this equation, Z_{wl} (= $R_{wl}(f)l_1 + j\omega L_{wl}l_1$) and Z_{w2} (= $R_{w2}(f)l_2 + j\omega L_{w2}l_2$) represent the impedances of the bonding wires used between the RF feeding line and EAM's metal pad (length: l_1) and between the EAM's pad and load impedance (length: l_2), respectively. Thus, these parameters consist of the frequency-dependent conduction resistance $R_{wl(2)}(f)$ and inductance $L_{wl(2)}$. Z_m denotes the impedance of EAM, R_m is the series resistance, and C_m is the junction capacitance. R_{pho} ($\cong dV_m/dl_{ph}$) is the equivalent AC resistance caused by the dependence of the photocurrent I_{ph} (governed by electro-absorption) at the modulator junction voltage V_m [10]. This optical power dependent current path results in the drop of V_m . As the frequency increases, its influence (i.e. the drop of V_m) is reduced due to the decrease of impedance governed by C_m (i.e. $1/\omega C_m$). As a result, it flattens the frequency responses of EML. If the EAM is properly designed, the parasitic capacitance C_p should be very small, while the leakage resistance R_j should be much larger than $1/\omega C_m$. On the other hand, in order to eliminate the power consumption of the load Z_L (or the Z_L and ground) and its influence on the RF frequency responses can be neglected as long as $C_{cb} > 0.5$ pF.



Fig. 1. (a) Equivalent circuit diagram of EML with input sources and load, (b) equivalent circuit of EAM. Z_{in} represents the input impedance $(=Z_{wl}+1/(1/Z_m+1/(Z_{w2}+Z_L)))$ and the bold line indicates the RF feeding line.

To evaluate the effects of adjusting the lengths of bonding wires, l_1 and l_2 , on the frequency response of EML, we analyzed the RF gain G(f) (= $20log_{10}[V_m(f)/V_m(0)]$) and return loss R(f) (= $20log_{10}[(Z_{in} - R_S)/(Z_{in} + R_S)]$) by using the circuit parameters listed in Table 1. In this analysis, we extracted L_{wl} , L_{w2} , R_{mb} and C_m from the measured data using a bonding wire with a diameter of 25.4 µm and an EML with a 175-µm long EAM, respectively.

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Table 1. Device parameters

Parameter	Value	Description
R_S	50Ω	Source impedance
Z_L	50 Ω	Load impedance
$R_{wl(2)}$	$\Omega \cdot mm^{-1} \cdot GHz^{-1/2}$	Conduction resistance of bonding wire
$L_{wl(2)}$	0.5 nH/mm	Inductance of bonding wire
R_m	5 Ω	Series resistance of EAM
C_m	0.42 pF	Junction capacitance of EAM
R_{pho}	125 Ω	Photocurrent-induced equivalent AC resistance of EAM
C_p	$10^{-3} \mathrm{pF}$	Parasitic capacitance of EAM
$\hat{R_j}$	$10^8 \Omega mm$	Leakage resistance of EAM

Figure 2 shows the frequency-dependent RF gain G(f) and the return loss R(f) calculated as a function of l_1 . In this calculation, l_2 is assumed to be zero. The result shows that, as l_1 increases, the bandwidth of G(f) decreases and the R(f) improves at lower frequencies. This is because the increase of Z_{w1} accelerates the drop of V_m (mainly by $\omega L_{w1}l_1$) and prevents the decrease of Z_{in} resulting from C_m (mainly by $R_{w1}(f)l_1$). Thus, in order to prevent the bandwidth reduction caused by the wire inductance between the RF feeding line and EML chip, it is needed to reduce L_{w1} as much as possible by shortening l_1 (i.e. $\leq 300 \ \mu$ m), using a bonding wire with a large diameter, or utilizing multiple bonding wires. The result in Fig. 2 also shows that, if the EML package is implemented with only one boning wire (i.e., $L_{w2} = 0 \ nH/mm$) [2,3], its modulation bandwidth cannot be enhanced to be larger than that of the EML chip itself (i.e., $L_{w1} = L_{w2} = 0 \ nH/mm$).

Figure 3 shows that, unlike the effect of l_1 , the bandwidth of G(f) increases and the R(f) improves at lower frequencies as l_2 increases. It should be noted that the modulation bandwidth of EML is increased mainly due to the RF resonance caused by C_m and L_{w2} . However, the bandwidth of G(f) does not increase any more due to the overshoot in G(f) when $l_2 > 1$ mm. Therefore, to increase the bandwidth of G(f) and improve the R(f), we need L_{w2} must be as large as possible by increasing l_2 under no excessive overshoot in G(f).



Fig. 2. (a) Frequency-dependent RF gain G(f) and (b) RF return loss R(f) for various values of l_1 . In this calculation, l_2 is assumed to be zero.

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Fig. 3. (a) Frequency-dependent RF gain G(f) and (b) RF return loss R(f) for various values of l_2 . In this calculation, l_1 is assumed to be zero.

The effect of bonding wire on the modulation bandwidth of EML can be different for different value of C_m . Figure 4 shows the 3-dB bandwidth of the RF gain G(f) and the peak value of G(f) (i.e., overshoot) calculated as a function of l_2 for various values of C_m . In this calculation, l_1 is assumed to be 300 µm. This figure shows that the 3-dB bandwidth of G(f) is increased and then gradually saturated as l_2 increase, and its saturation length is getting shorter as the value of C_m is smaller.



Fig. 4. (a) 3 dB-bandwidth of the RF gain G(f) and (b) peak value of G(f) calculated as a function of l_2 for various values of C_m . In these calculations, l_1 is assumed to be 300 µm.

3. Experimental results

Figure 5 shows a photograph of the EML packaged on the sub-mount, which contains a load resistor (nominal 50 Ω) and an RF feeding line (made of the flexible printed circuit board (FPCB)) in the form of the grounded coplanar waveguide (GCPW). The load resistor and *RF* feeding line were die bonded on the sub-mount by using the silver-paste at the temperature of 150 °C. By using a network analyzer, we measured the S-parameters of this package as follows. S11 and S21 for the feeding line were measured to be <-21.5 dB and >-0.2 dB @ 20 GHz, respectively, and S11 for the load resistor was <-17 dB @ 20 GHz. We used a lumped EML chip, developed for 10-Gb/s operation [2]. The fiber-coupled output power with the EAM biased 0 V and the capacitance including the electrode pad of EAM were ~3 dBm@70

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mA and ~0.42 pF@-1.2 V, respectively. Under this condition, the 3-dB modulation bandwidth of this EML chip was measured to be ~16 GHz.



Fig. 5. A photograph of the EML on the sub-mount.

We prepared two types of bonding wires in the EML package: for the type 1, $l_1 \cong l_2 \cong 300$ µm, and, for the type 2, $l_1 \cong 300$ µm and $l_2 \cong 700$ µm. Both of these bonding wires had a diameter of 25.4 µm, and were wedge-bonded on the metallic pads. Figure 6 shows the E/O responses and return losses measured by using these two types of bonding wires. The fast decays near 3 GHz in the E/O response due to the carrier accumulation in the isolation region between the DFB LD and EAM) [11]. The result clearly showed that the modulation bandwidth of EML was increased as we increased l_2 . In fact, it showed that the 3-dB bandwidth was increased from ~18 GHz to ~28 GHz. This improvement was achieved by the *RF* resonance caused by the junction capacitance of EAM and the inductance of bonding wire. From these results, we concluded that, by using the proposed technique, it is possible to utilize the EML chip previously developed for 10-Gb/s operation in the 25-Gb/s and 40-Gb/s systems. This technique is cost-effective since it utilizes the conventional packaging process.



Fig. 6. Measured E/O responses and return losses of the EML packaged in type 1 ($l_1 \cong l_2 \cong 300 \mu$ m, dashed lines) and type 2 ($l_1 \cong 300 \mu$ m, $l_2 \cong 700$ um, solid lines) measured at the laser current of 70 mA and the bias voltage of -1.2V.

4. Summary

We proposed and demonstrated a simple and cost-effective technique for improving the modulation bandwidth of EML by using the RF resonance caused by the EML chip and bonding wire. By using an EAM circuit model, we analyzed the effects of adjusting the lengths of the bonding wires, l_1 and l_2 , on the frequency responses of EML. To verify this analysis, we packaged a lumped EML chip (having a modulation bandwidth of ~16 *GHz*) on the sub-mount and measured its frequency responses. The measured data showed the same tendency with the simulated results. The measured data also showed that we could increase the modulation bandwidth of this EML to 28 GHz simply by adjusting the lengths of the bonding wires to be $l_1 \cong 300 \,\mu\text{m}$, $l_2 \cong 700 \,\mu\text{m}$. From these results, we concluded that the proposed technique could be useful for the development of the cost-effective high-speed light sources.

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